# Dynamic Analysis of Gough-Stewart Platform Manipulator 

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#### Abstract

A novel derivation to evaluate all the controlled forces which cause by the motors and effected along the prismatic joints on the legs of the Gough-Stewart platform manipulator based on the virtual work method is proposed in this paper. In this paper the manipulator can be considered as a multibody mechanism with rigid elements. It can be assumed that the manipulator motion was known. The aim of the dynamic analysis in this paper is to evaluate all the controlled forces which necessary to implement the manipulator programming motion.


Keywords: Gough-Stewart, Robotics, Dynamic Analysis and Manipulator
الحسابات الايناميكية للمناول نوع كوف - ستيوارت

تم في هذا البحث استخدام طريقة جديدة واشتقاق جديد لحساب كل القوى المسيطر عليها التي تولدها محركات الروبوت في المفاصل الطولية في ارجل الروبوت نوع ستيو ارت . لقد تم اعتماد طريقة احنساب الجهـ الافتر اضي بعد ان تم اعتبار الروبوت الية ميكانيكية متكونة من عدة اجسام صلبة, وان حركة الروبوت معلومة ما بين نقطتين معلومتين . لقد تم اعتبار الارجل بدون اية انحناءات جانبية ناتجة من جراء تاثير القوى المختلفة المؤثرة عليها. ان الهرف الرئيسي من هذا البحث هو ايجاد طريقة علمية لاحتساب جميع القوى المسيطر عليها في محركات الروبوت والتي تعنبر ضرورية لانجاز الحركة المبرمجة للروبوت بين معلومتين .

نقطنين

## 1. Introduction.

In general, Gough-Stewart platform manipulator is a six degree of freedom with two main bodies [3]. The fixed body is called the base, while another body is regarded as movable and is called the moving plate. These two bodies are connected together by six extensible legs. In this paper we assumed that every leg of the legs of the manipulator is consists of two parts connected together with a prismatic kinematic joint (p). The prismatic joints are affected under the controlled forces which cause by the motors. All the legs are connected with the base by spherical kinematic joints (s) in the points $A_{i}$, while they are connected with the moving plate by spherical joints with fingers in the points $B_{i}$, as shown in Figure1.

Force and moment analysis for any robotic system are useful to choose the suitable motors for implementation the programming motion [5]. Several methods are proposed to derive dynamic equations of the GoughStewart manipulator.

The motivation of this paper is to derive a mathematical formulation for evaluation all the controlled forces in the prismatic kinematic joints of the Gough-Stewart platform manipulator.

## 2. Position equation

In order to describe the motion of the moving plate of the manipulator relative to the base, we assumed that the moving plate attached to the base as shown in Figure 1, so the position
vector of the point $A$ in the Global coordinate system can be written as:

$$
r_{A i}=\left(x_{A i}, y_{A i}, z_{A i}\right)^{T}, i=1,2, . ., n
$$

and the position vector of the point $B$ in the local coordinates system can be written as:

$$
r_{B i}=\left(x_{B i}, y_{B i}, z_{B i}\right)^{T}, i=1,2, \ldots, n .
$$

The rotation of transformation between the local (moving plate) and the global coordinate systems can be represented by Euler's rotation matrix with $\psi$ as pitch angle, $\theta$ as yaw angle and $\varphi$ as a roll angle [1]. Thus $\mathbf{R}_{\mathbf{0 1}}$ is a matrix of coordinate transformation from the local to the global coordinates can be written as:

$$
\mathbf{R}_{01}=\left(\begin{array}{lll}
A_{11} & A_{21} & A_{31} \\
A_{12} & A_{22} & A_{32} \\
A_{13} & A_{23} & A_{33}
\end{array}\right)
$$

where

$$
\begin{aligned}
& A_{11}=\cos (\psi) \cos (\varphi)-\sin (\psi) \cos (\theta) \sin (\varphi), \\
& A_{12}=\sin (\psi) \cos (\varphi)+\cos (\psi) \cos (\theta) \sin (\varphi), \\
& A_{13}=\sin (\theta) \sin (\varphi), \\
& A_{21}=-\cos (\psi) \sin (\varphi)-\sin (\psi) \cos (\theta) \cos (\varphi), \\
& A_{22}=-\sin (\psi) \sin (\varphi)+\cos (\psi) \cos (\theta) \cos (\varphi), \\
& A_{23}=\sin (\theta) \cos (\varphi), A_{31}=\sin (\psi) \sin (\theta), \\
& A_{32}=-\cos (\psi) \sin (\theta), \operatorname{and} A_{33}=\cos (\theta) .
\end{aligned}
$$

Generally, kinematic relations of manipulator are expressed by the loop equations. In Figure 1 one of the loops are represented, and its equation can be written as a leg vector:
$q_{i}=l_{i}=r_{0}+R_{01} r_{B i}-r_{A i} ;$
$i=1,2, \ldots, n$
where $r_{0}$ the position vector of the moving plate centre in the Global coordinate system,
Thus:
$q_{i}{ }^{2}=r_{0}{ }^{T} r_{0}+r_{B i}{ }^{T} r_{B i}+r_{A i}{ }^{T} r_{A i}+2 r_{0}{ }^{T} R_{01} r_{B i}$

$$
-2 r_{A i}^{T} R_{01} r_{B i}-2 r_{0}^{T} r_{A i}
$$

$$
i=1,2, \ldots, n
$$

$\Psi_{\Pi i}\left(q_{i}, x, y, z, \psi, \theta, \varphi\right)=r_{0}{ }^{T} r_{0}+r_{B i}{ }^{T} r_{B i}+r_{A i}{ }^{T} r_{A i}$
$+2 r_{0}^{T} R_{01} r_{B i}-2 r_{A i}^{T} R_{01} r_{B i}-2 r_{0}^{T} r_{A i}-q_{i}^{2}=0 ;$
$i=1, .6 . n$

It is assumed that, the geometrical parameters of the moving plate centre can be represented as:

$$
\rho=\left(\begin{array}{llllll}
x & y & z & \psi & \theta & \varphi
\end{array}\right)^{T},
$$

and the legs vectors as:

$$
q=\left(\begin{array}{lllll}
q_{1} & q_{2} & . & . & q_{6}
\end{array}\right)^{T} .
$$

Thus the position equation will be written as:

$$
\begin{aligned}
& \Psi_{\Pi}(q, \rho)= \\
& {\left[\Psi_{\Pi 1}\left(q_{1}, x, . ., \varphi\right) \quad \Psi_{\Pi 2}\left(q_{2}, x, . ., \varphi\right) \quad \mathrm{K} \quad \Psi_{\Pi 6}\left(q_{6}, x, . ., \varphi\right)\right]^{T}}
\end{aligned}
$$

$$
\begin{equation*}
\Psi_{\Pi}(q, \rho)=0 \tag{1}
\end{equation*}
$$

## 3. Forces and moments effected on the moving plate of the manipulator. <br> 3.1. Gravity force.

It can be considered that $\bar{G}_{1}$ is the gravity force vector of the moving plate in the Global coordinate system, so its projector on the Global coordinate is $G_{1}^{(0)}=\left(\begin{array}{lll}0 & 0 & -m_{1} g\end{array}\right)^{T}$ and his projector on the local coordinate system is $G_{1}^{(1)}=R_{10} G_{1}^{(0)}$.
There is a moment that caused by the gravity force, its vector relative to the beginning of the local coordinate system can be given as:

$$
\begin{equation*}
\bar{M}_{01}\left\{G_{1}\right\}=\bar{r}_{C 1} \times \bar{G}_{1} \tag{2}
\end{equation*}
$$

### 3.2. Working force.

It can be assumed that the moving plate affected under a working force with a vector $\bar{P}_{H}$. This force will be caused a moment with a vector relative to the beginning of the local coordinate system $\bar{M}_{01}\left\{\bar{P}_{H}\right\}$.

### 3.3. Controlled forces(motors forces).

Every leg of the legs of the manipulator affected under a controlled force which effected along the prismatic joint (along the leg). This force causes by the motor $Q_{i}$

## 4. Inertia forces and moments evaluation.

The formulation of the inertia force of the moving plate of the manipulator can be written as:

$$
\begin{equation*}
\Phi_{1}=-m_{1} W_{C 1} \tag{3}
\end{equation*}
$$

where $W_{C 1}$ is the vector projector of the acceleration of the moving plate centre on the local coordinate system [7]. Inertia moment of the moving plate relative to the beginning of the local coordinate system can be obtained as:

$$
\begin{equation*}
\bar{L}_{01}=-\overline{\bar{J}}_{C 1} \cdot \bar{\varepsilon}_{1}-\bar{\omega}_{1} \times \overline{\bar{J}}_{C 1} \cdot \bar{\omega}_{1}+\bar{r}_{C 1} \times \bar{\Phi}_{1} \tag{4}
\end{equation*}
$$

In the above equation $\overline{\bar{J}}_{C 1}$ - inertia tenser of the moving plate relative to its mass centre; $\bar{\omega}_{1}$ - moving plate angular velocity; $\bar{\varepsilon}_{1}$ - moving plate angular acceleration and $\bar{r}_{C 1}$ - radius vector of the moving plate centre. Thus the projector of the inertia moment on the local coordinate system can be written as:

$$
\begin{equation*}
L_{01}=-J_{C 1} \cdot \varepsilon_{1}-\omega_{1} J_{C 1} \omega_{1}+r_{C 1} \times \Phi_{1} \tag{5}
\end{equation*}
$$

where

$$
\begin{aligned}
& r_{C 1}=\left(\begin{array}{ccc}
0 & -z_{C 1} & y_{C 1} \\
z_{C 1} & 0 & -x_{C 1} \\
-y_{C 1} & x_{C 1} & 0
\end{array}\right), \\
& \omega_{1}=\left(\begin{array}{ccc}
0 & -\omega_{z} & \omega_{y} \\
\omega_{z} & 0 & \omega_{x} \\
-\omega_{y} & \omega_{x} & 0
\end{array}\right) .
\end{aligned}
$$

Similarly the inertia forces and moments for the legs of the manipulator can be written as:

$$
\Phi_{S}=-m_{S} W_{C S}
$$

and,

$$
\bar{L}_{0 S}=-\overline{\bar{J}}_{C S} \cdot \bar{\varepsilon}_{S}-\bar{\omega}_{S} \times \overline{\bar{J}}_{C S} \cdot \bar{\omega}_{S}+\bar{r}_{C S} \times \bar{\Phi}_{S}
$$

But we assumed that the mass, inertia forces and moments of the legs is very little, so we will neglect them in the evaluations of the present work.

## 5. Virtual work evaluation.

The moving plate of the manipulator affected under the total force which consists of its weight, working force and inertia force. The total force vector is

$$
\begin{equation*}
\overline{P_{1}}=\bar{G}_{1}+\bar{P}_{H}+\bar{\Phi}_{1} . \tag{6}
\end{equation*}
$$

The total force causes a total moment relative to beginning of the local coordinate system as follow:

$$
\begin{equation*}
\bar{M}_{01}=\bar{M}_{01}\left\{\bar{G}_{1}\right\}+\bar{M}_{01}\left\{\bar{P}_{H}\right\}+\bar{L}_{01} \tag{7}
\end{equation*}
$$

Thus the summation of the virtual work of the moving plate by any little transformation and orientation of the moving plate can be obtained as:
$\delta A_{1}=\bar{P}_{1} \cdot \delta \bar{r}_{C 1}+\bar{M}_{01} \cdot \delta \bar{\gamma}_{1}$
In the above equation $\delta \bar{r}_{c 1}-$ a moving plate centre little translation, $\delta \bar{\gamma}_{1}-\mathrm{a}$ moving plate little rotations and it can be obtained as follow:

From Figure 2 it can be assumed that the moving plate angular orientation relative to Global coordinate by using Euler's angles $\delta \psi_{1}, \delta \theta_{1}, \delta \varphi_{1}$. The angular orientation can be written as:

$$
\begin{align*}
\delta \bar{\gamma}_{1} & =\delta \bar{\psi}_{1}+\delta \bar{\theta}_{1}+\delta \bar{\varphi}_{1} \\
& =\bar{k}_{0} \delta \bar{\psi}_{1}+\bar{n} \delta \theta_{1}+\bar{k}_{1} \delta \varphi_{1}, \tag{9}
\end{align*}
$$

Equation (9) can be written in a following form:

$$
\begin{equation*}
\delta \gamma_{1}=\Gamma_{1} \delta \Delta_{1} \tag{10}
\end{equation*}
$$

In the above equation $\delta \Delta_{1}=\left(\begin{array}{l}\delta \psi_{1} \\ \delta \theta_{1} \\ \delta \varphi_{1}\end{array}\right)$,
$\Gamma_{1}=\left(\begin{array}{ccc}\sin \varphi_{1} \sin \theta_{1} & \cos \varphi_{1} & 0 \\ \sin \theta_{1} \cos \varphi_{1} & -\sin \varphi_{1} & 0 \\ \cos \theta_{1} & 0 & 1\end{array}\right)$.
Thus:
$\delta \bar{\gamma}_{1}=\bar{i}\left(\delta \theta_{1} \cos \varphi_{1}+\delta \psi_{1} \sin \varphi_{1} \sin \theta_{1}\right)$
$+\bar{j}_{1}\left(-\delta \theta_{1} \sin \varphi_{1}+\delta \psi_{1} \sin \theta_{1} \cos \varphi_{1}\right)$
$+\bar{k}_{1}\left(\delta \psi_{1} \cos \theta_{1}+\delta \varphi_{1}\right)$

Now, the virtual work can be derived in a new relation as follow:
$\delta A_{1}=P_{1}^{T} \delta r_{C 1}+M_{01}^{T} \Gamma_{1} \delta \Delta_{1}$
If we assume that all the forces and moments effected on the moving plate can be written as:
$F_{1}=\left(\begin{array}{llllll}P_{x 1} & P_{y 1} & P_{z 1} & M_{01 x} & M_{01 y} & M_{01 z}\end{array}\right)^{T}$,
$K_{1}=\left(\begin{array}{cc}E_{3} & 0 \\ 0 & \Gamma_{1}\end{array}\right)$
and the geometrical parameters of the moving plate centre can be represented as:

$$
\rho=\left(\begin{array}{llllll}
x_{\mathrm{cl}} & y_{\mathrm{cl}} & z_{\mathrm{cl}} & \psi_{1} & \theta_{1} & \varphi_{1}
\end{array}\right)^{T}
$$

the virtual work will be written as:

$$
\begin{equation*}
\delta A_{1}=F_{1}^{T} K_{1} \delta \rho \tag{12}
\end{equation*}
$$

## 6. The controlled forces evaluation.

In the static equilibrium condition the virtual work represented in equation (12) equal to virtual work which causes by the controlled forces which can given in the following form [9]:

$$
\begin{equation*}
\delta A_{1}=Q^{T} \delta q \tag{13}
\end{equation*}
$$

Thus:

$$
\begin{equation*}
Q=\left(\frac{\delta \rho}{\delta q}\right)^{T} K_{1}^{T} F_{1} \tag{14}
\end{equation*}
$$

From equation (1) we can derive the following relation:

$$
\begin{array}{r}
\frac{\partial \Psi_{\Pi}}{\partial q} \delta q+\frac{\partial \Psi_{\Pi}}{\partial \rho} \delta \rho=0 \\
\text { and } \frac{\delta \rho}{\delta q}=-\left(\frac{\partial \Psi_{\Pi}}{\partial \rho}\right)^{-1}\left(\frac{\partial \Psi_{\Pi}}{\partial q}\right) .
\end{array}
$$

Where

$$
\begin{aligned}
\frac{\partial \Psi_{\Pi}}{\partial q} & =\left(\begin{array}{ccc}
\frac{\partial \Psi_{\Pi 1}}{\partial q_{1}} & \ldots & \frac{\partial \Psi_{\Pi 1}}{\partial q_{6}} \\
\ldots & \ldots & \ldots \\
\frac{\partial \Psi_{\Pi 6}}{\partial q_{1}} & \ldots & \frac{\partial \Psi_{\Pi 6}}{\partial q_{6}}
\end{array}\right) \\
& =\operatorname{diag}\left\{2 q_{i}\right\}
\end{aligned}
$$

and,

$$
\frac{\partial \Psi_{\Pi}}{\partial \rho}=\left(\begin{array}{ccc}
\frac{\partial \Psi_{\Pi 1}}{\partial x_{1}} & \ldots & \frac{\partial \Psi_{\Pi 1}}{\partial \varphi_{1}} \\
\ldots & \ldots & \ldots \\
\frac{\partial \Psi_{\Pi 6}}{\partial x_{1}} & \ldots & \frac{\partial \Psi_{\Pi 6}}{\partial \varphi_{1}}
\end{array}\right)
$$

Thus the controlled forces (14) can be formulated as follows:

$$
\begin{equation*}
Q=\left[\left(\frac{\partial \Psi_{\Pi}}{\partial \rho}\right)^{-1}\left(\frac{\partial \Psi_{\Pi}}{\partial q}\right)\right]^{T} K_{1}^{T} F \tag{15}
\end{equation*}
$$

From the above equation can evaluate all the controlled forces which effected on the prismatic joints of the manipulator.

## 7. Example of analysis:

In this example a manipulator with the following parameters is considered:
a. Coordinates of points $A_{i}$ of the base (in the Global coordinates system) are as in the matrix:
$r_{A i}=\left(\begin{array}{ccccccc}2 & 1.4140 & -1.414 & -2 & -1.414 & 0 & 1.414 \\ 0 & 1.4142 & 1.414 & 0 & -1.414-2 & -1.414 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0\end{array}\right)$,
b. Coordinates of points $B_{i}$ in the platform coordinates systems are as in the matrix:
$r_{B i}=\left(\begin{array}{cccccccc}1.6 & 1.45 & -0.278-0.917-1.575 & -1.31 & 0.278 & 0.917 \\ 0 & 0.676 & 1.575 & 1.31 & -0.278-0.917-1.575-1.31 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0\end{array}\right.$,
$c$. It can be assumed that the platform
moves from the point 1 to the point 6 throw the track points $2,3,4,5$. The coordinates of the center of the platform $\mathrm{P}\left({ }^{r_{0}}\right)$ and Euler angles of its orientation $(\psi, \theta, \varphi)$ are as shown in Table 1 .
d. At any point of the track it can be found that:
$q_{i}=r_{0}+A_{0,1} r_{B i}-r_{A i} ; i=1,2, \ldots, 6$
The results for all points are as explained in Table 2.
$e$. The platform mass is $\mathrm{m}=15 \mathrm{Kg}$, tenser of platform Inertia is

$$
J=\left(\begin{array}{ccc}
0.1 & 0 & 0 \\
0 & 0.1 & 0 \\
0 & 0 & 0.1
\end{array}\right) \mathrm{Kg} \mathrm{~m}^{2}
$$

and the forces and moments vector which applied to the platform is:
$P=\left(\begin{array}{llllll}0.5 \mathrm{~N} & 0.25 \mathrm{~N} & -0.75 \mathrm{~N} & 0.5 \mathrm{Nm} & 1 \mathrm{~N} . \mathrm{m} & 1.25 \mathrm{~N} . \mathrm{m}\end{array}\right)^{T}$
$f$. The movement of the platform throw the track points (1,2,3,4,5,6) with acceleration and angular velocity of point P as are follow:

$$
\Leftrightarrow\left[\begin{array}{c}
0.1 \mathrm{~m} / \mathrm{s}^{2} \\
0.01 \mathrm{~m} / \mathrm{s}^{2} \\
0.002 \mathrm{~m} / \mathrm{s}^{2} \\
0.15 \mathrm{deg} / \mathrm{s}^{2} \\
0.1 \mathrm{deg} / \mathrm{s}^{2} \\
0.01 \mathrm{deg} / \mathrm{s}^{2}
\end{array}\right]
$$

$$
\text { And } \quad \omega=\left[\begin{array}{c}
0.8 \mathrm{deg} / s \\
0.16 \mathrm{deg} / s \\
0.01 \mathrm{deg} / s
\end{array}\right]
$$

.The results of the motorized forces are as shown in the Table 3. In this table it can be found that there is a singularity position in the point 4 , it can be seen that, the forces are very high in this position.

## 8. Conclusions

In this paper, an approach for the dynamic analysis of the Gough-Stewart platform manipulator is proposed based on the virtual work. The mathematical simulation for the robot arms in this paper is a novel method and it has been derive by the authors. The stiffness of the arms caused by the motion of the platform was neglected. The joints of the manipulator's arms have been assumed as ideal joints (without friction forces). It has been proved that all the controlled forces can be evaluated without using the reactions in the kinematic joints

## 9. References

[1] Kolovsky M.Z., Evgrafov A.N., Semenov Y.A., Sloush A.V. "Advanced theory of mechanisms and machines". Translated by Lilov L., Springer, 2000.
[2] Luc Baron, and Jorge Angeles. "The direct kinematics of parallel manipulators under joint - sensor redundancy". IEEE, Transaction on Robotics and Automation, February 2000.vol.16, No.1, c.12_19.
[3] Min _ Jie Liu, Cong _ Xin Li, and Chong _Ni Li. "Dynamics analysis
of the Gough _Stewart platform manipulators". IEEE, Transaction on Robotics and Automation, February 2000.vol.16, No.1, c.94_98.
[4] Litvin F. L., Zhang Yi, Parenti Castelli V., Innocenti C."Singularities, configurations, and displacement functions for manipulators". Int. J. Robotics res. 1986. V. 5.No.2. p. 52-65.
[5] Miomir Vukobratovic, Veljko Potkonjak. "Dynamics of manipulation robots theory and application". Springer - Verlag Berlin, Heidelberg. 1982.303p.
[6] Nair R. Maddocks J.H. "On the forward kinematics of the parallel manipulators". The International Journal of Robotics Research, Vol 13, No.2, April 1994, pp. 171 - 188.
[7] McCarthy J. M. "Dual orthogonal matrices in manipulator kinematics". Int. J. Robotics res. 1986. V. 5.No.2. p. 45-51.
[8] Thomas Geike, John P. "Inverse dynamic analysis of parallel manipulators with full mobility . Int. J. Robotics res. 2006.
[9] Nicolas Andreff. "Dynamic analysis and simulation of a six degree of freedom Stewart platform manipulator ". The international journal of robotic researches, Vol. 26 , No. 7 . 2007, p 677-687.
[10] S. Kemal Ider. "Inverse dynamics of parallel manipulators in the presence of drive singularitieß . Journal of mechanical engineering science, Vol. 220 , No.1. 2006 , p 61-72.
[11]. Alexie Sokolov, Paul Xirouchakis. "Dynamics analysis of a 3-DOF parallel manipulator with

R-P-S joint structure". Journal of mechanisms and machine theory, Vol. 42, Issue 5, 2007. p 541-557.

Table (1) Platform center coordinates and its angles of orientation

| pooint <br> coord. | 1 | 2 | 3 | 4 | 5 | 6 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{X}(\mathrm{m})$ | 0.25 | 0.5 | 0.85 | 1 | 1.2 | 1.4 |
| $\mathrm{Y}(\mathrm{m})$ | -0.2 | -0.45 | -0.8 | -1 | -1.3 | -1.5 |
| $\mathrm{Z}(\mathrm{m})$ | 2.2 | 2.35 | 2.7 | 3 | 3.2 | 3.5 |
| $\psi^{\circ}$ | -10 | -20 | -30 | -30 | -40 | -45 |
| $\theta^{\circ}$ | 12 | 18 | 35 | 30 | 35 | 40 |
| $\varphi^{\circ}$ | 8 | 15 | 26 | 31.95832 | 35 | 40 |

Table (2) Legs length

| point | 1 | 2 | 3 | 4 | 5 | 6 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| q1 | 2.2664 | 2.5530 | 3.2884 | 3.6198 | 4.1003 | 4.6117 |
| q2 | 2.6035 | 3.0504 | 3.9605 | 4.2101 | 4.7743 | 5.2914 |
| q3 | 2.5986 | 2.9654 | 3.7631 | 3.9677 | 4.3756 | 4.7763 |
| q 4 | 2.5846 | 2.9322 | 3.5634 | 3.7983 | 4.1697 | 4.5303 |
| q 5 | 2.2393 | 2.3956 | 2.6738 | 3.1315 | 3.3431 | 3.6687 |
| q 6 | 2.0335 | 2.0585 | 2.1544 | 2.6063 | 2.7686 | 3.0855 |

Table (3) Controlled forces

| point <br> Force <br> In (N) | 1 | 2 | 3 | 4 | 5 | 6 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{Q}_{1}$ | -1.832 | -2.760 | 12.306 | 4.58 E 7 | 10.507 | 5.258 |
| $\mathrm{Q}_{2}$ | 0.057 | 1.082 | -13.461 | -4.26 E 7 | -12.351 | -7.239 |
| $\mathrm{Q}_{3}$ | 3.943 | 4.478 | -4.062 | -2.49 E 7 | -0.546 | 2.234 |
| $\mathrm{Q}_{4}$ | -3.822 | -4.945 | 9.763 | 3.85 E 7 | 4.538 | -0.069 |
| $\mathrm{Q}_{5}$ | -0.867 | -0.874 | 3.758 | 2.34 E 7 | 6.022 | 3.860 |
| $\mathrm{Q}_{6}$ | 1.640 | 2.133 | -9.906 | -4.09 E 7 | -9.394 | -5.173 |
| $\sum \mathrm{Q}^{2}\left(N^{2}\right)$ | 36.974 | 58.619 | 556.73 | 8.25 E 15 | 408.367 | 126.723 |



Figure (1) Position of the robot joints


Figure (2) The rotation of the axis

